

INTRODUCTION TO

FAST DYNAMOS

[AND SLOW DYNAMOS.]

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UNIVERSITY OF EXETER

14:00 - 15:00

MONDAY.

OBJECTIVES:

- UNDERSTAND BASIC DYNAMO MECHANISMS IN FLUID FLOWS (CLASSIFICATION),
- OBTAIN ASYMPTOTIC / APPROXIMATE GROWTH RATES,
- UNDERSTAND BASIC SATURATION MECHANISM.

STARTING POINT:

- INDUCTION EQUATION, NON-DIMENSIONALISED USING SCALE \mathcal{L} AND MAGNITUDE \mathcal{U} OF THE FLUID FLOW:

$$\partial_t \underline{B} = \nabla \times (\underline{u} \times \underline{B}) + \varepsilon \nabla^2 \underline{B}$$

$$\nabla \cdot \underline{B} = 0$$

- $\varepsilon \equiv R_m^{-1}$, $R_m = \frac{\mathcal{U} \mathcal{L}}{\eta}$ MAGNETIC REYNOLDS NUMBER
- KINEMATIC THEORY: \underline{u} GIVEN FLUID FLOW, (NOT RANDOM OR TURBULENT).

$$\partial_t \underline{B} = \nabla \times (\underline{u} \times \underline{B}) + \epsilon \nabla^2 \underline{B}$$

OR

$$\underbrace{\partial_t \underline{B} + \underline{u} \cdot \nabla \underline{B}}_{\text{ADVECTION}} = \underbrace{\underline{B} \cdot \nabla \underline{u}}_{\text{STRETCHING}} + \underbrace{\epsilon \nabla^2 \underline{B}}_{\text{DIFFUSION.}}$$

WITH $\nabla \cdot \underline{B} = 0, \quad \nabla \cdot \underline{u} = 0.$

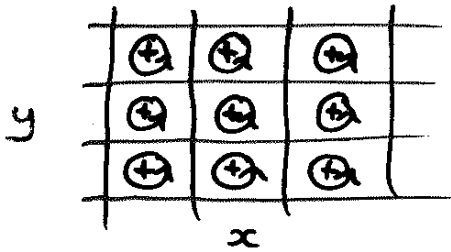
FOCUS ON THE LIMIT $R_m \rightarrow \infty$ OR $\epsilon \rightarrow 0,$

i.e. OF WEAK DIFFUSIVITY (WITH FIXED CHOICE OF FLOW FIELD \underline{u}).

WHY?

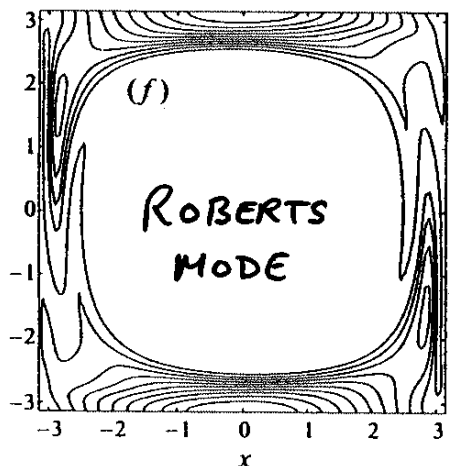
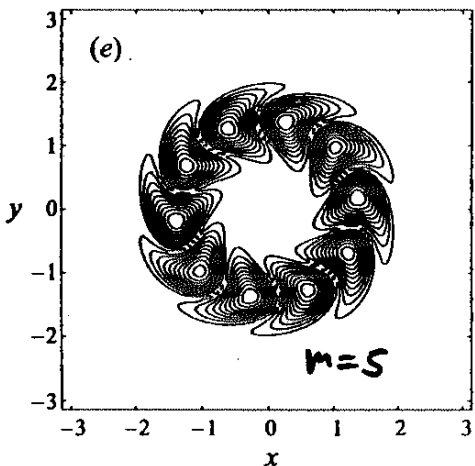
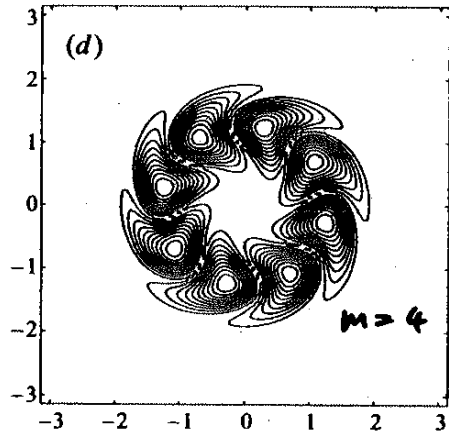
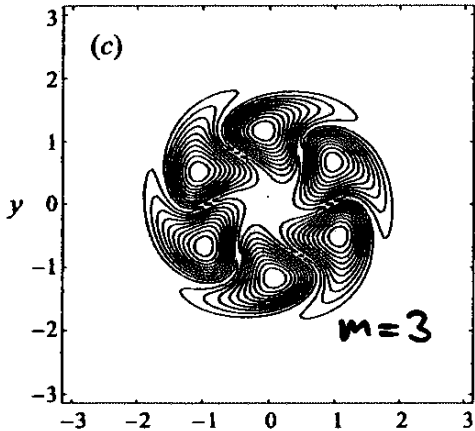
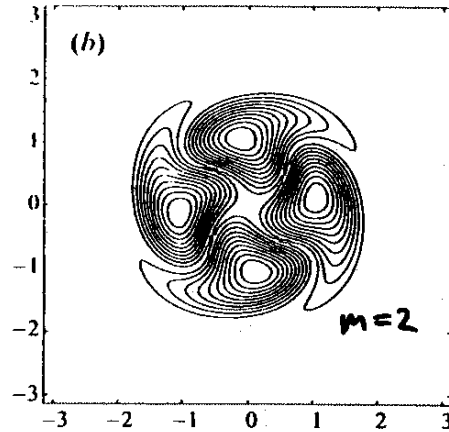
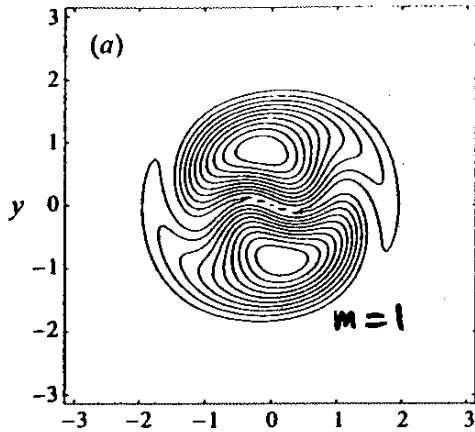
- ASTROPHYSICALLY IMPORTANT LIMIT (STARS, GALAXIES)
- ALLOWS ASYMPTOTIC APPROXIMATION GIVING USEFUL FORMULAE & INFORMATION.
- ALLOWS CLASSIFICATION OF DYNAMO MECHANISMS.

$$\underline{u} = (v_y, -v_x, kv), \quad \Psi = a(1 + \cos x)(1 + \cos y).$$



"FAST BREEDER REACTOR"

MAGNETIC FIELD IN 1 CELL OF FLOW:



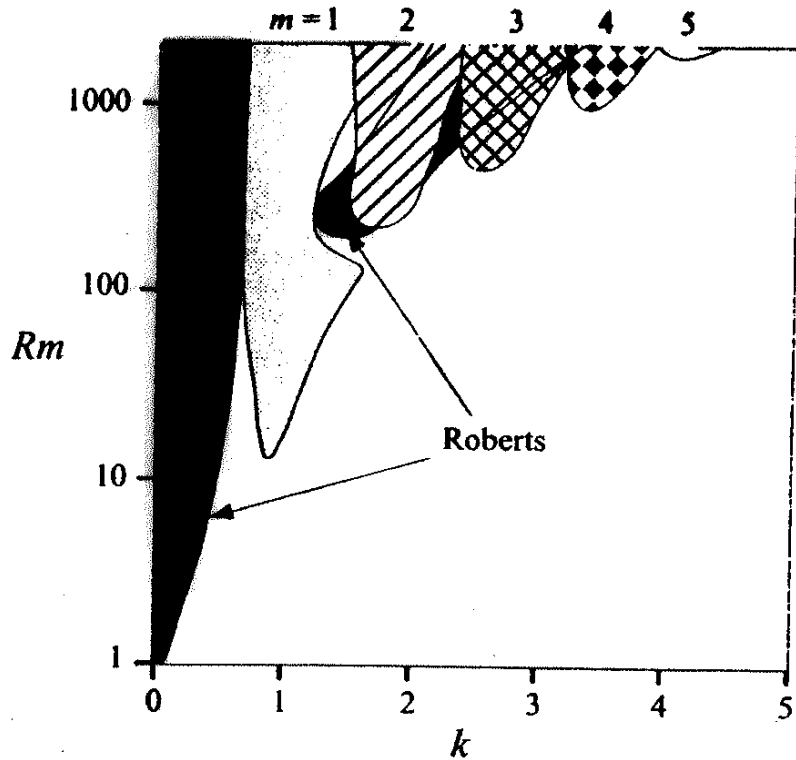
PONOMARENKOV MODES

NUMERICAL RESULTS:

$$\underline{B} \propto e^{ikz + pt}$$

$$\gamma = \text{Re } p$$

CRITICAL
VALUES OF
 R_m (FOR
ZERO GROWTH
RATE, $\gamma=0$).

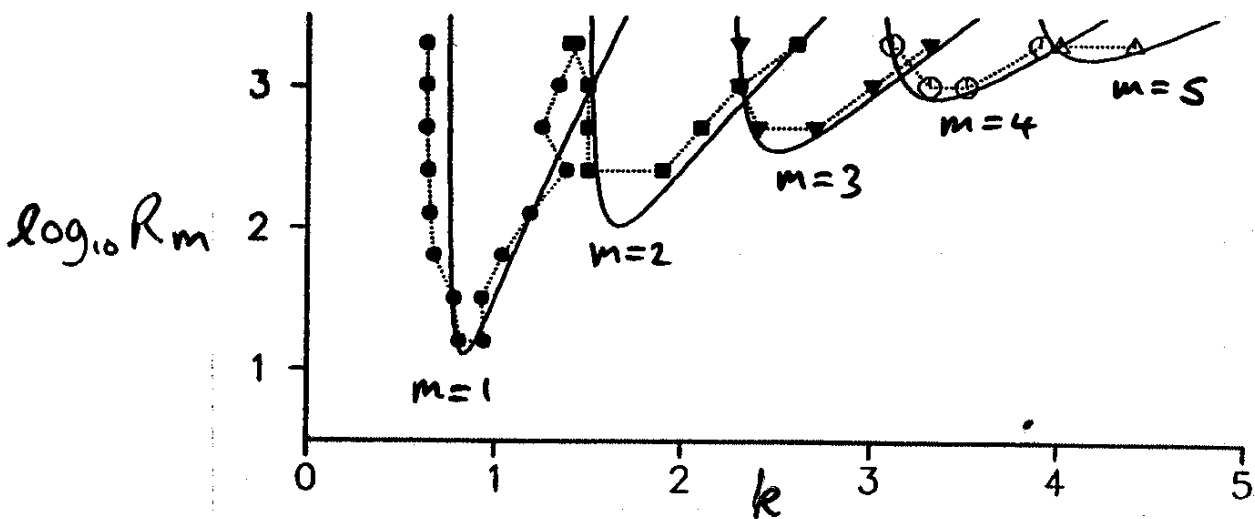


ASYMPTOTIC (LARGE- R_m) RESULTS



GILBERT,
PONTY, GAFD
93, 2000.

NUMERICAL RESULTS ● ■ ▼ ○ △



EXCELLENT AGREEMENT BETWEEN HIGH- R_m THEORY

& CRITICAL R_m VALUES.

BASIC MECHANISMS :

PONOMARENKO DYNAMO (SMOOTH):

Ponomarenko
1973

$$\underline{u} = r\Omega(r)\underline{e}_\theta + U(r)\underline{e}_z$$

$$\underline{B} = e^{im\theta + ikz + pt} \underline{b}(r)$$



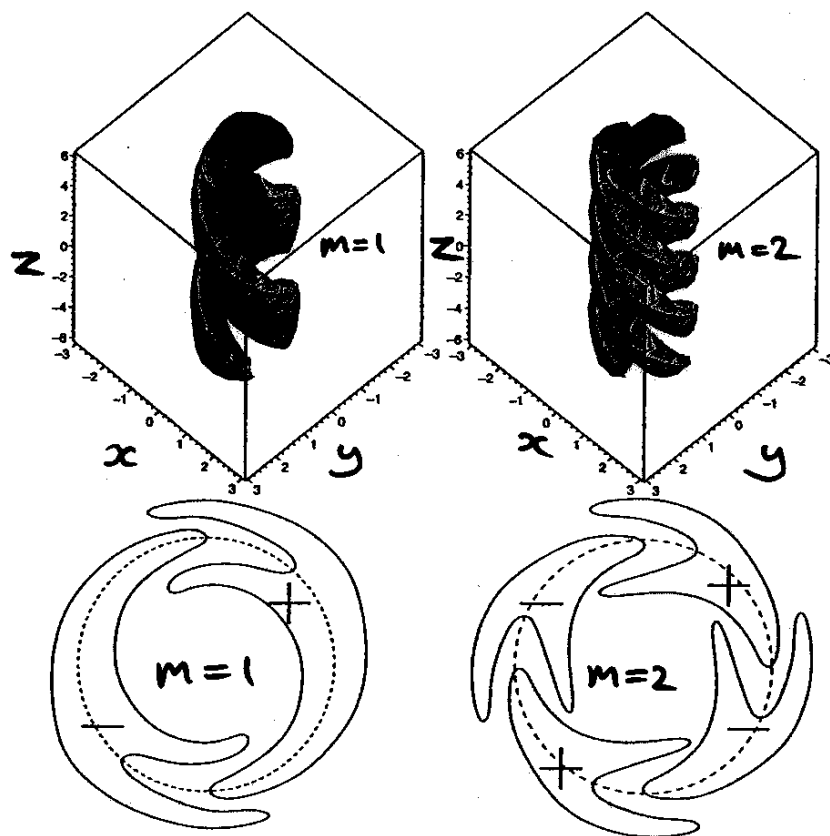
Gailitis,
Freiberg

+
z-motion

m, k ARE WAVENUMBERS WE CAN CHOOSE.

AS $\epsilon \rightarrow 0$ FIELD LOCALISES AT A RADIUS a WHERE

$$m\Omega'(a) + kU'(a) = 0$$



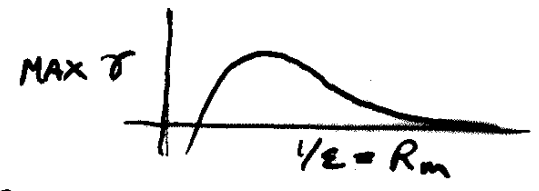
$$\gamma = \text{Re } p = \sqrt{\epsilon |m\Omega'(a)|/a} - \frac{1}{2} \sqrt{\epsilon |m\Omega''(a) + kU''(a)|} \\ - \epsilon (m^2/a^2 + k^2).$$

Ruzmaikin, Sokoloff, Shukurov, Gilbert 1988.

FASTEST GROWING MODES HAVE:

$$m, k = O(\epsilon^{-1/3})$$

$$\gamma = O(\epsilon^{1/3})$$



"SLOW" DYNAMO: MAXIMUM $\gamma \rightarrow 0$ AS $\epsilon \rightarrow 0$

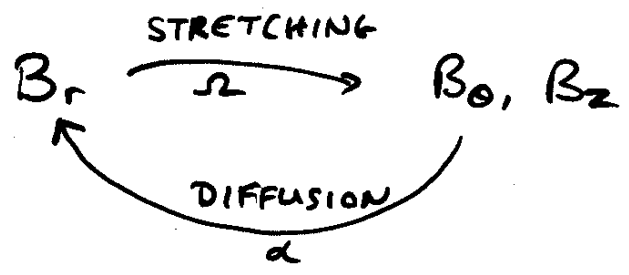
• DYNAMO ACTION (AT LARGE R_m) ONLY

POSSIBLE AT A GIVEN RADIUS IF:

$$r \left| \frac{\Omega''(r)}{\Omega'(r)} - \frac{U''(r)}{U'(r)} \right| < 4$$

(GEOMETRICAL CONDITION).

• MECHANISM:



DYNAMO IS

DIFFUSION-LIMITED.

• CAN GENERALISE TO FLOWS WITH

CLOSED STREAM SURFACES $\psi(x, y)$.



SIMILAR MECHANISM, GROWTH RATES,

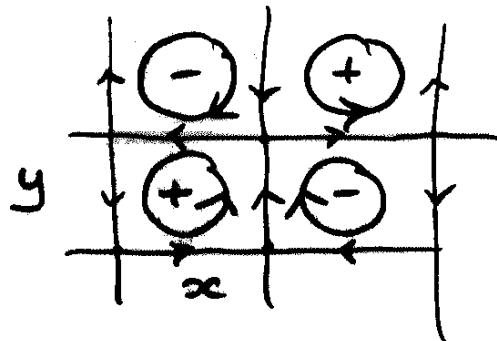
GEOMETRICAL CONDITION.

CELLULAR STEADY FLOWS:

G.O. Roberts 1970.

$$\underline{u} = (\sin x \cos y, -\cos x \sin y, K \sin x \sin y)$$

$$K = \sqrt{2}$$



BELTRAMI FLOW

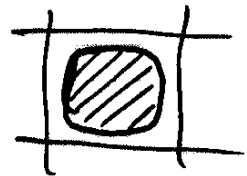
$$\nabla \times \underline{u} = K \underline{u}$$

ABC FLOW

$A=B, C=0$, RESCALE

Low R_m : CLASSICAL α^2 DYNAMO.

HIGH R_m : PONOMARENKO MODES



Soward
1990

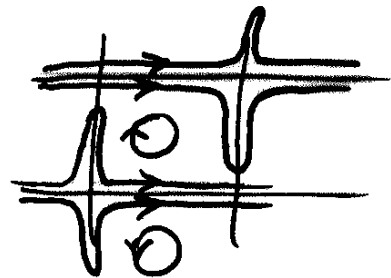
BUT ONLY WHERE GEOMETRICAL

CONDITION IS MET $////$

BUT NOW HAVE "ROBERTS" MODES

FIELD LOCALISED ON SEPARATRICES

& X-POINTS. "SFS" MECHANISM:



PUT $\underline{B} \propto e^{ikz + pt}$

$$\gamma = \alpha k - \varepsilon k^2, \quad \alpha = -\frac{1}{2} K \varepsilon^{1/2} G$$

$$G \approx 1.0655$$

Childress 1979

Soward 1986

VALID FOR $k = O(1)$.

FASTEST GROWING MODES HAVE:

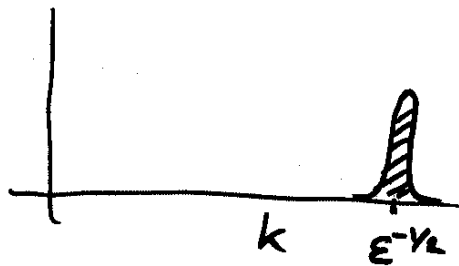
$$\gamma = O\left(\frac{\log \log \varepsilon^{-1}}{\log \varepsilon^{-1}}\right), \quad k = O\left(\frac{\varepsilon^{-1/2}}{\sqrt{\log \varepsilon^{-1}}}\right)$$

GROWTH RATE GOES TO ZERO AS $\varepsilon \rightarrow 0$

(BUT ONLY LOGARITHMICALLY) — SLOW DYNAMO.

FIELD HAS TO ADOPT SMALL SCALES AS

SPECTRUM $E_M(k)$



DIFFUSION STILL
LIMITING GROWTH.

NO LARGE-SCALE FIELD COMPONENT.

IMPORTANT SLOW DYNAMO MECHANISM.

ALSO CATS' EYE DYNAMOS.

PONOMARENKO



& SEPARATRIX MODES

PLANE LAYER

Ponty, Gilbert, Soward.

DEFINITIONS:

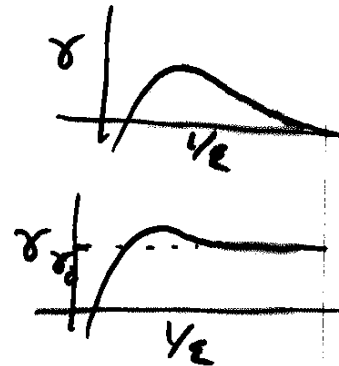
SOLVE THE INDUCTION EQUATION
FOR A GIVEN FLOW \underline{u} AND LET
 $\gamma(\epsilon)$ BE THE MAXIMUM GROWTH RATE,
MAXIMISED OVER ALL MODES / INITIAL CONDITIONS.

PUT:

$$\gamma_0 = \lim_{\epsilon \rightarrow 0} \gamma(\epsilon)$$

\underline{u} IS A SLOW DYNAMO IF $\gamma_0 \leq 0$

A FAST DYNAMO IF $\gamma_0 > 0$



Vainshtein +
Zeldovich 1972.

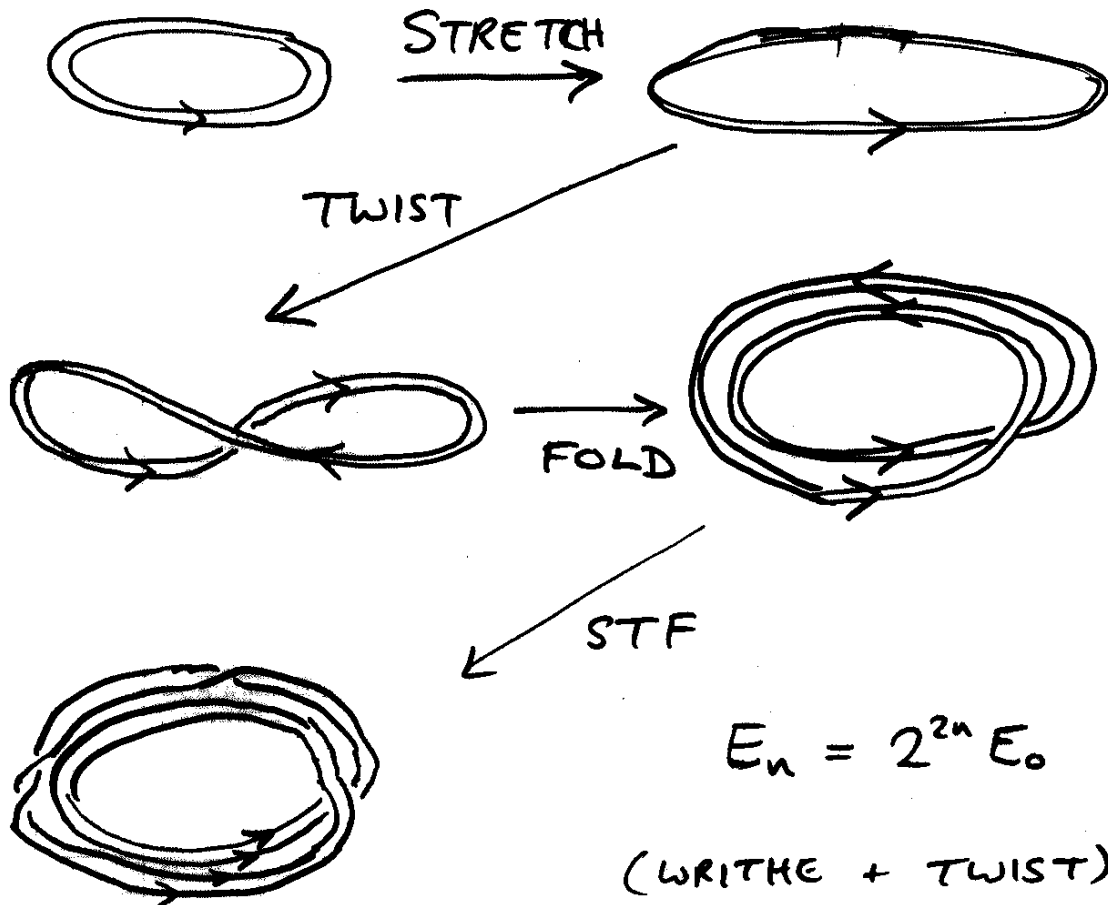
FAST DYNAMO:

NEED TO AVOID RELYING ON DIFFUSION
TO AMPLIFY A COHERENT FIELD.

STRETCH - TWIST - FOLD (STF)

MECHANISM / **PICTURE**

Vainshtein, Zeldovich 1978
Alfvén.



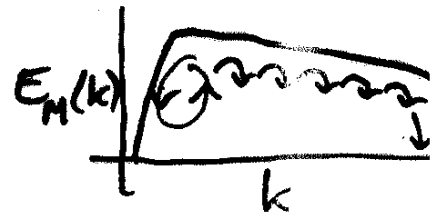
$$E_n = 2^{2n} E_0$$

(WRITHE + TWIST)

"PICTURE" OF FAST DYNAMO: $\gamma_0 = \log 2$

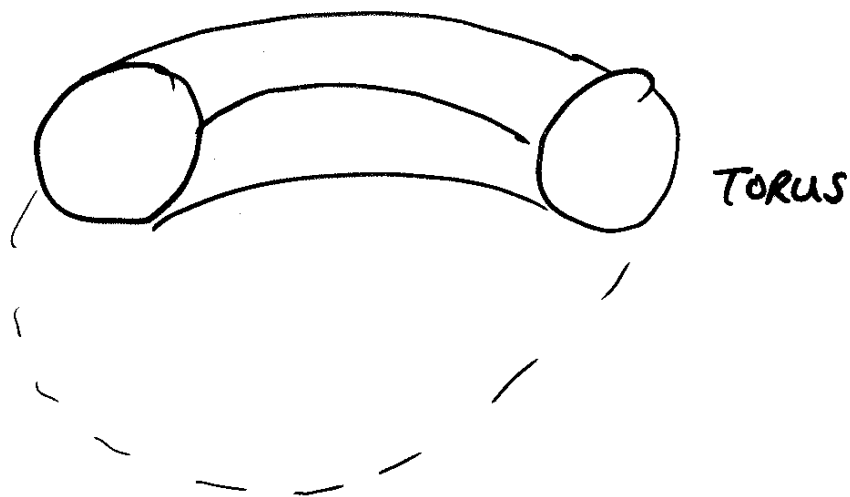
- CHAOTIC (LAGRANGIAN SENSE) FLOW
- AMPLIFIES LARGE-SCALE COHERENT FIELD

+ CASCADE OF FLUCTUATIONS

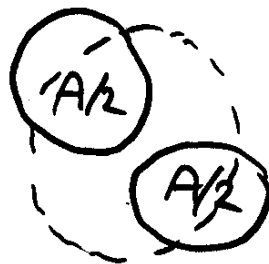
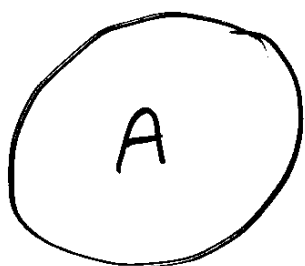


- PICTURE HARD TO REALISE

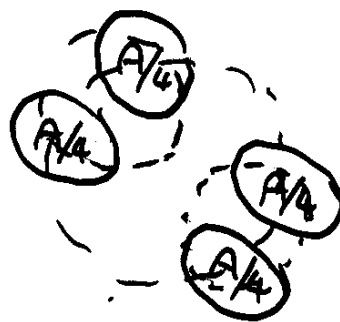
PROBLEM WITH S.T.F.



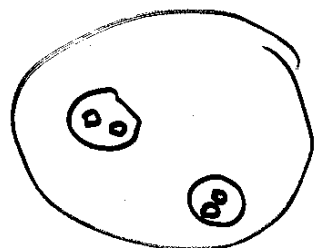
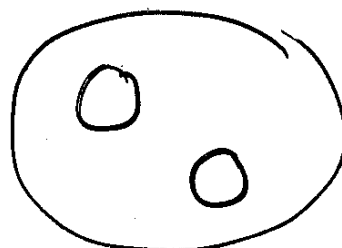
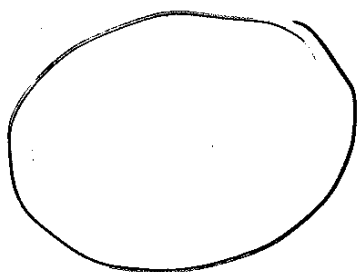
VOLUME-PRESERVING
MAP/FLOW



FIELD MUST
ESCAPE FROM TORUS.



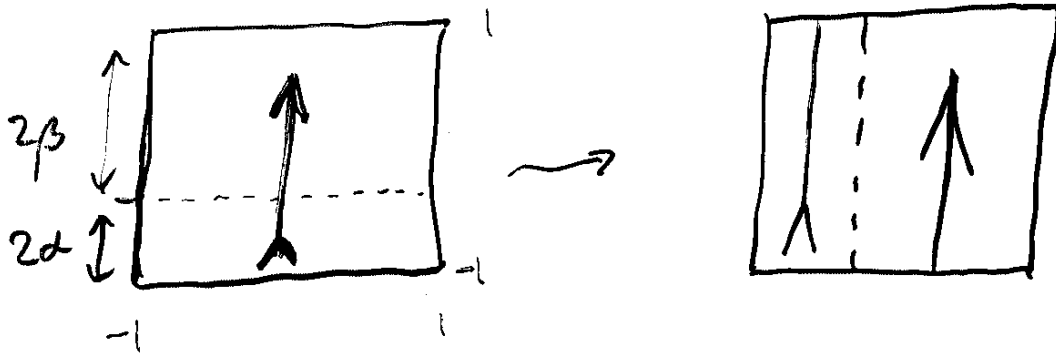
FOR A VOLUME-CONTRACTING FLOW: $\text{div } \underline{u} < 0$:



"SOLENOID MAP"

BAKERS MAP MODEL OF STF

Film / Ok.

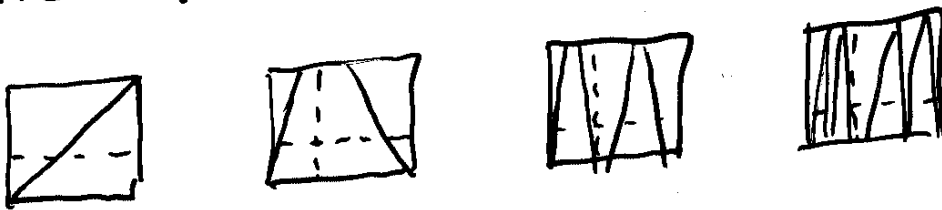


$\alpha + \beta = 1$

stretch by factor α^{-1} or β^{-1}

LAPUNOV EXPONENT $\lambda = \alpha \log \alpha^{-1} + \beta \log \beta^{-1}$

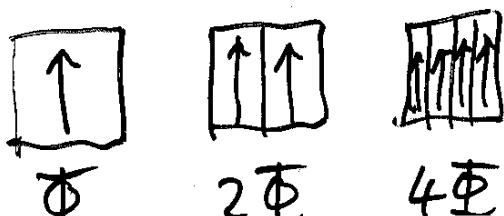
RATE OF GROWTH OF LINE LENGTH:



TOPOLOGICAL ENTROPY = LINE STRETCHING EXPONENT (HERE) = $h = \log 2 \geq \lambda$.

[DIFFERENT AVERAGING: LINES ACCUMULATE IN REGIONS OF STRONG STRETCHING.]

MAGNETIC FLUX (NO DIFFUSION)



$\gamma = \log 2 = h$
(NO FOLDING).

• RIGOROUS RESULT :

NEED LAGRANGIAN CHAOS (POSITIVE
 TOPOLOGICAL ENTROPY) FOR FAST DYNAMO
 ACTION IN A SMOOTH FLOW: $\gamma_0 \leq h$.

Klapper, Young. 1995
 (Also Finn, Ott)

• STEADY, 3-D ABC FLOWS

COMPUTATIONALLY INTENSIVE - PICTURE

Galloway, Frisch 1986

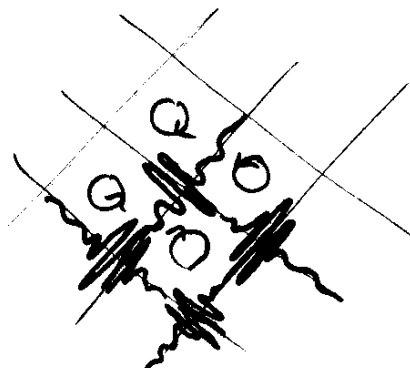
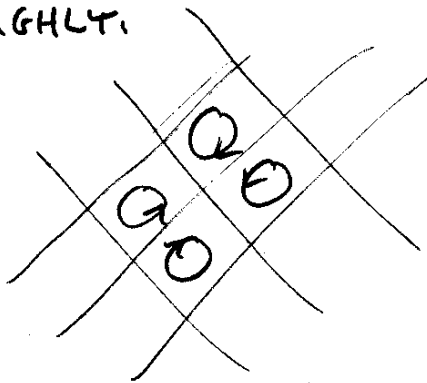
STILL UNCLEAR.

• UNSTEADY, 2-D FLOWS

Galloway, Proctor 1992
 Otani 1993

$$\underline{u} = 2 \cos^2 t (0, \sin x, \cos x) + 2 \sin^2 t (\sin y, 0, -\cos y)$$

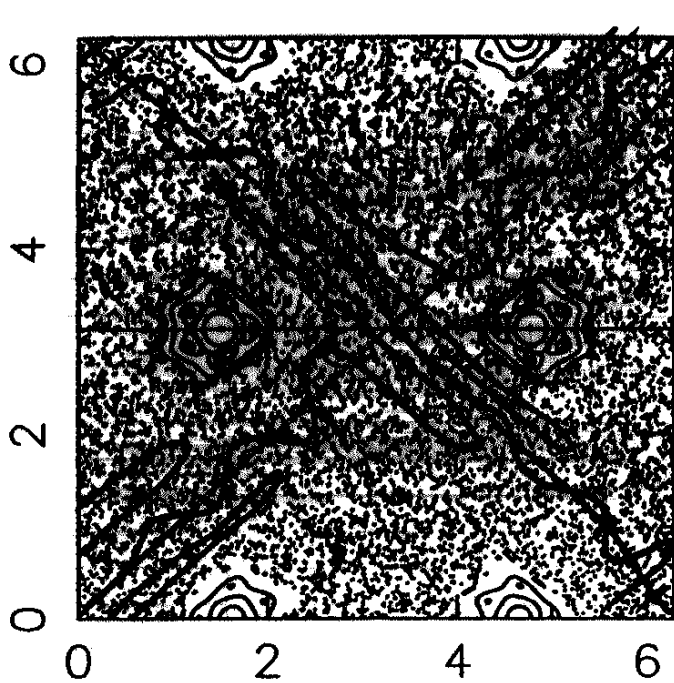
ROUGHLY:



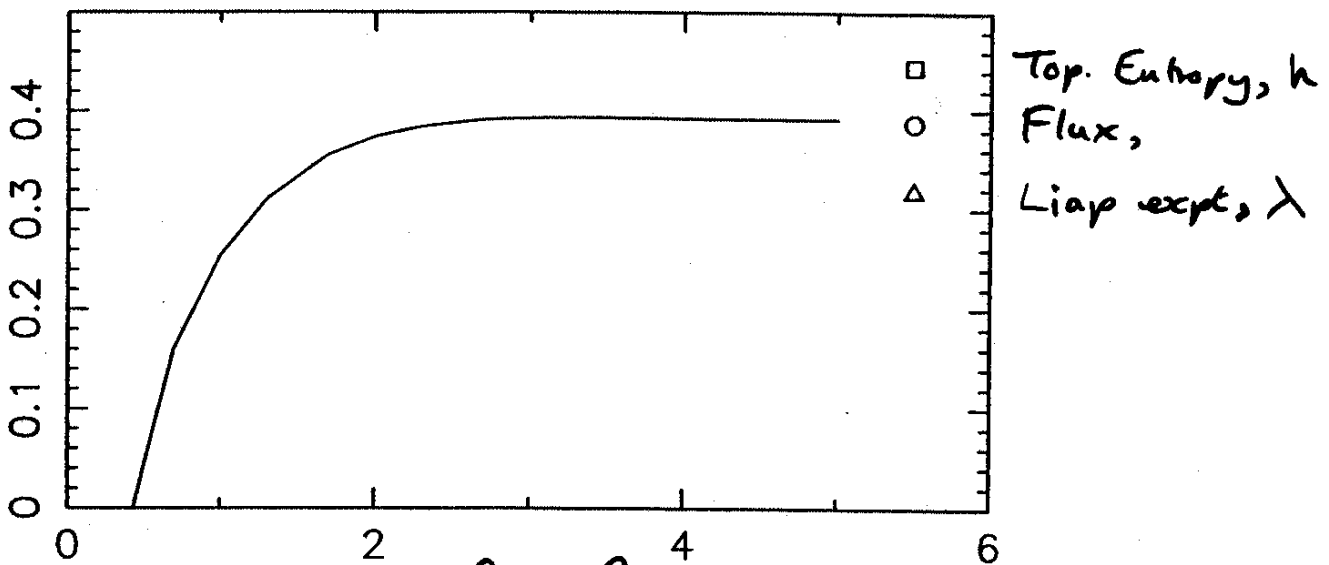
BROKEN
 SEPARATRICES
 -
 CHAOTIC
 LAYERS

ROBERTS FLOW + TIME DEPENDENCE.

POINCARÉ (STROBOSCOPIC) MAP, PERIOD 2π IN t .



FAST DYNAMO $\gamma_0 \approx 0.39,$



$$\log_{10} R_m = \log_{10} \frac{1}{\epsilon}$$

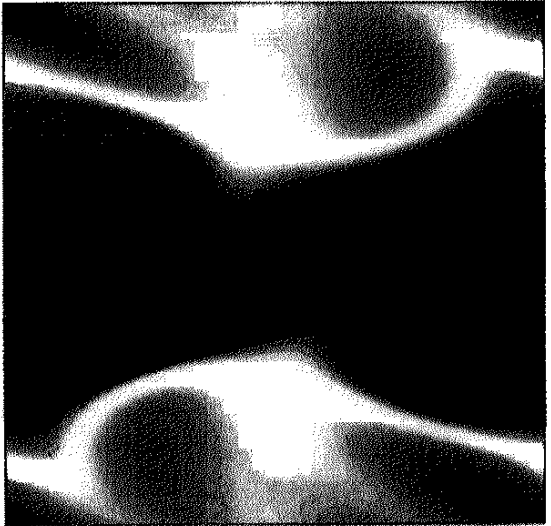
$$\underline{B} \propto e^{ikz + pt},$$

$$k = 0.8$$

LARGE-SCALE
FIELD COMPONENT

$$R_m = \frac{1}{\varepsilon}$$

$R_m = 10$



100

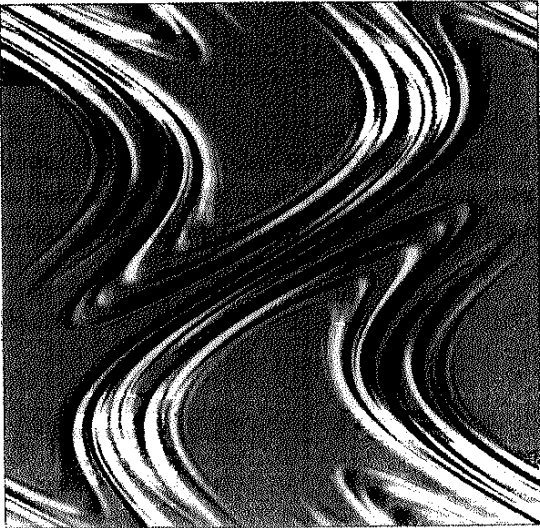
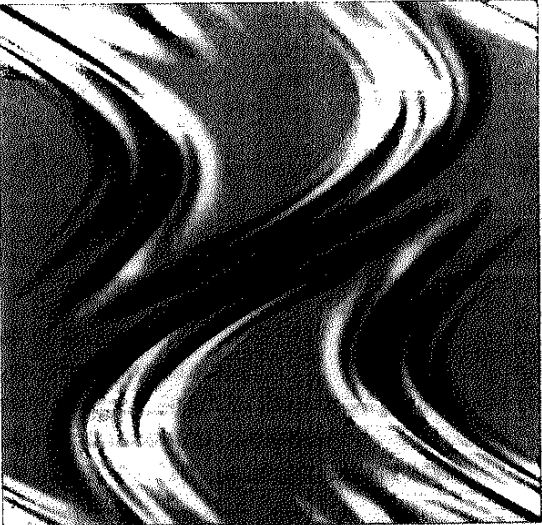
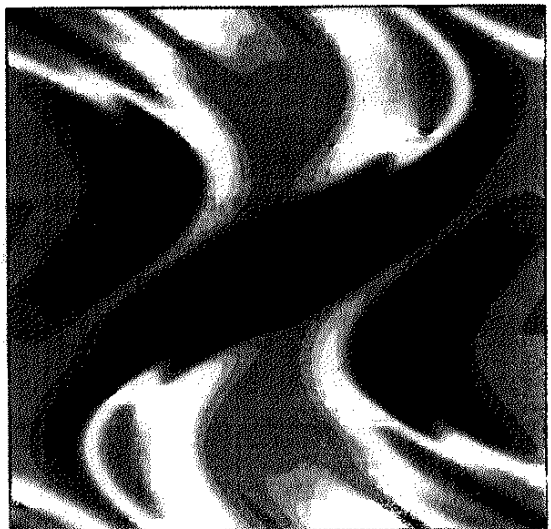


Fig. 2.3 Eigenfunctions for $k = 0.8$ and $\varepsilon = (a) 10^{-1}$, (b) 10^{-2} , (c) 10^{-3} , and 10^{-4} .

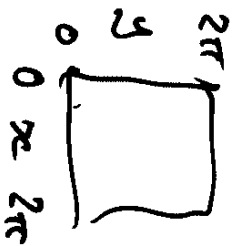
Ottavio's flow
eigenfunctions

$$k = 0.8$$

$$\varepsilon =$$

$$10^{-1} \quad 10^{-2}$$

$$10^{-3} \quad 10^{-4}$$



$$z \in [0, 2\pi]$$

$$x \in [0, 2\pi]$$

22

$R_m = 10^4$

SFS
 "STRETCH -
 FOLD -
 SHEAR"

MECHANISM

Bogly, Chilikess, Soward, Roberts, ...

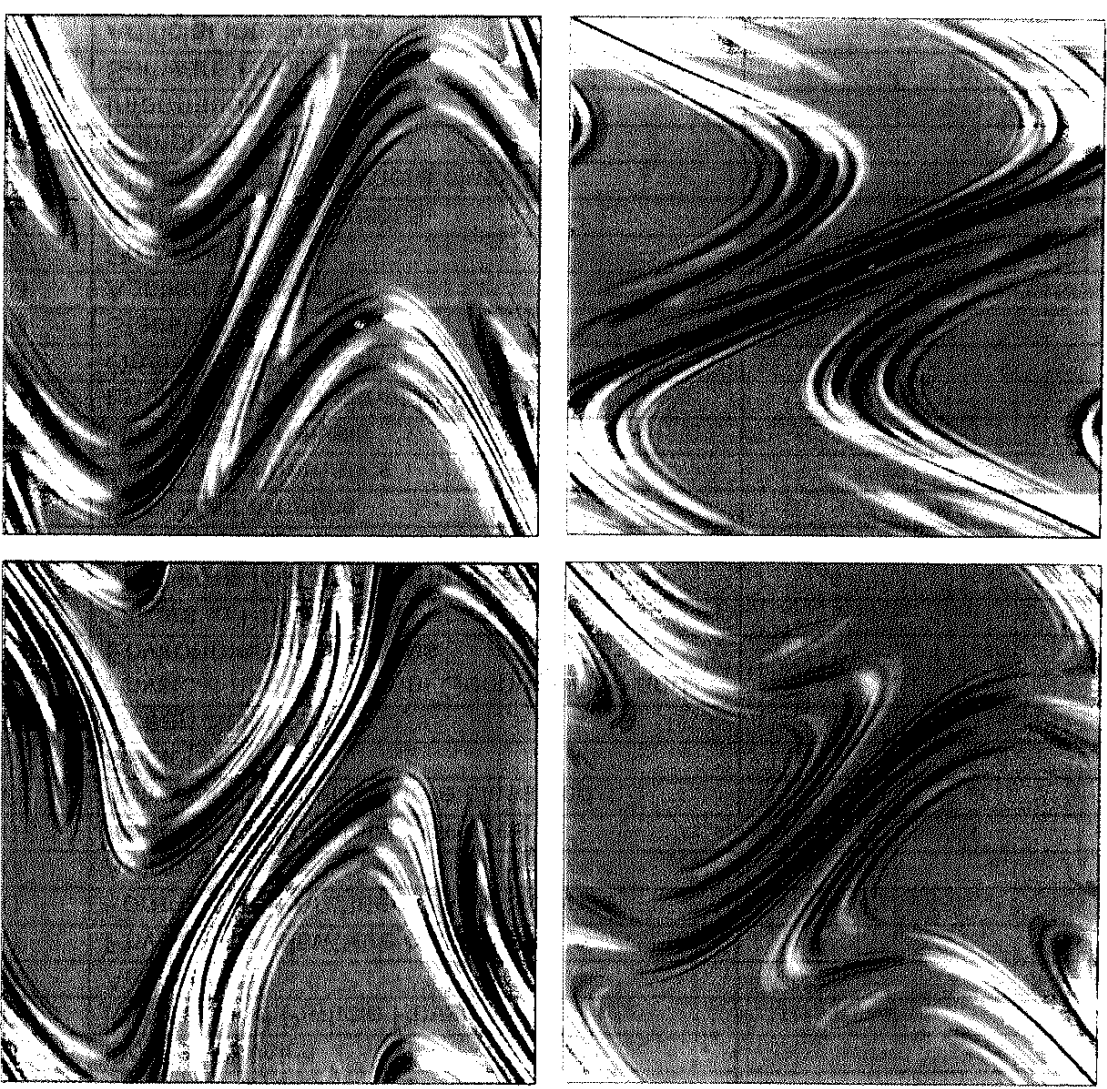
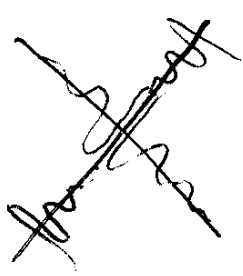


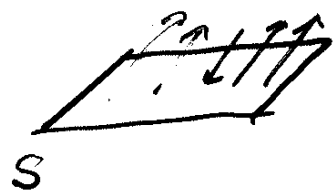
Fig. 2.5 Field evolution with $k = 0.8$ and $\epsilon = 10^{-4}$ for $t = 40\pi = (a) \pi/5, (b) 2\pi/5,$

Ostrov's flow
 $\epsilon = 10^{-4}$
 time - evolution
 over 1 cycle



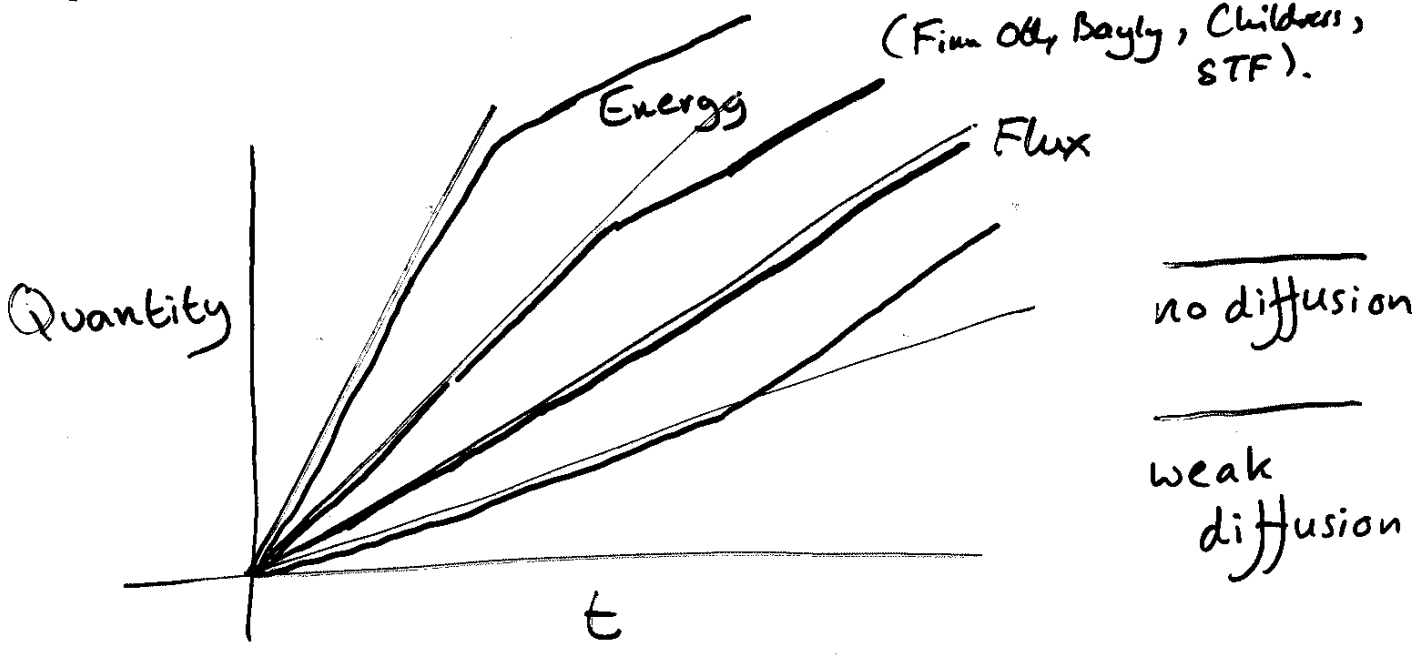
mechanism is piling
 up of sheets of
 field because
 of broken separatrix
 (Cf. Melnikov)
 + shear in Z-direction
 $\underline{u} = 2 \cos^2 t (0, \sin x, \cos x)$
 + $2 \sin^2 t (\sin y, 0, -\cos y)$.

MEASURE FLUX GROWTH:

$$F_s = \int_S \underline{B} \cdot d\underline{S} \propto e^{\Gamma_s t}$$


(NO DIFFUSION). CONJECTURE: $\Gamma_s = \gamma_0$

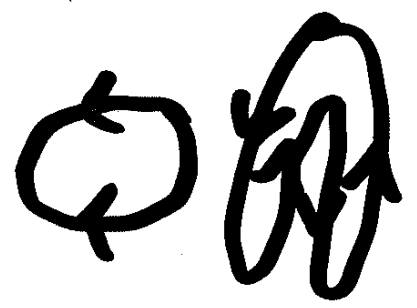
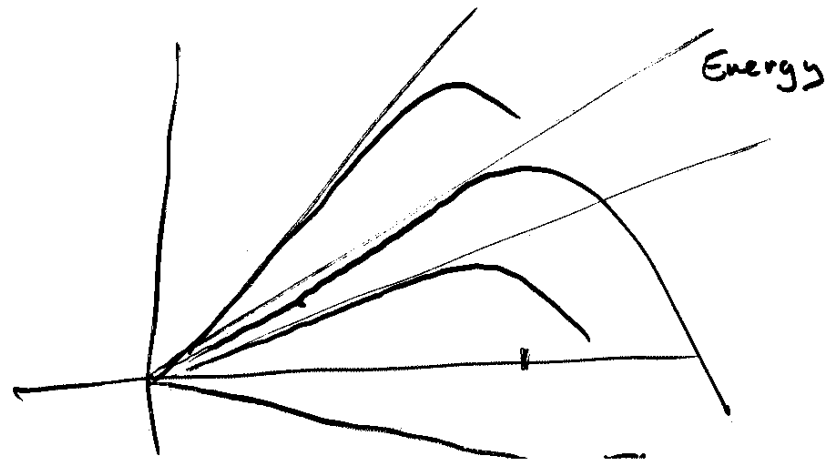
(From Old, Bayly, Childress, STF).



I.E. FLUX IS ONLY ROBUST LINK BETWEEN ZERO DIFFUSION AND WEAK DIFFUSION

[N.B. $\underline{u} = (u_1(x,y,t), u_2(x,y,t), 0)$

NOT A DYNAMO.



ANALYSIS: STUDY DYNAMOS IN MAPS.

APPLY MAP $M: b \rightarrow Tb$

APPLY WEAK DIFFUSION: $b \rightarrow H_\epsilon b$

"DYNAMO OPERATOR": $T_\epsilon \equiv H_\epsilon T$

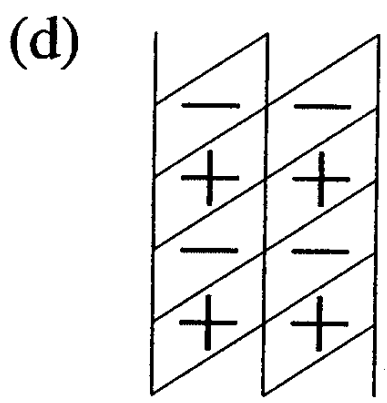
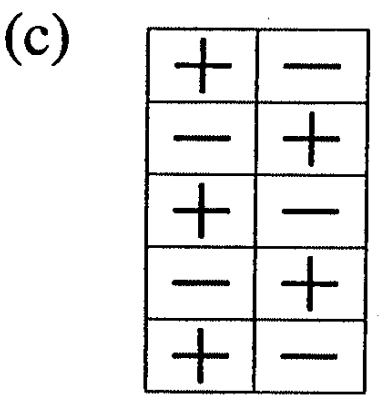
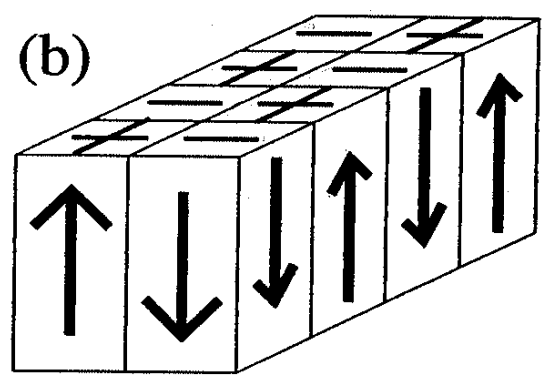
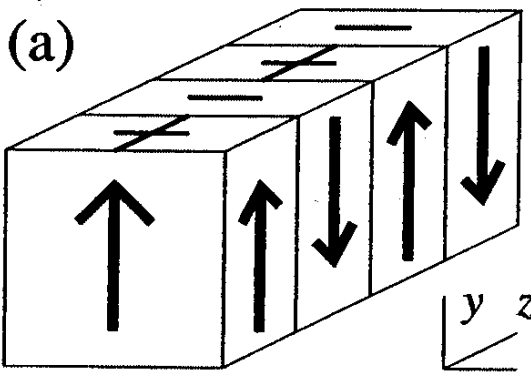
STUDY EIGENVALUES & EIGENFUNCTIONS OF

T_ϵ AS $\epsilon \rightarrow 0$. DYNAMO IS FAST IF

T_ϵ HAS EIGENVALUE(S) WITH $|\lambda|$ BOUNDED ABOVE 1 AS $\epsilon \rightarrow 0$

STRETCH-FOLD-SHEAR MECHANISM / MAP

Bayly, Childers



α : SHEAR PARAMETER.

START WITH:

$$\underline{B} = e^{ikz} b(x) \underline{e}_y + \text{complex conjugate.}$$

TAKE $k=1$ (WLOG).

DYNAMO OPERATOR (NO DIFFUSION) IS:

$$T b(x) = \begin{cases} 2 e^{-iax} b(1+2x) & -1 \leq x < 0. \\ -2 e^{-iax} b(1-2x) & 1 \geq x \geq 0 \end{cases}$$

FOLD GIVES
A MINUS
SIGN

STRETCHING
BY FACTOR
2

CONTRACTION OF
SCALE IN
x-DIRECTION.

SHEAR GIVES
A PHASE SHIFT

DIFFUSION: $\partial_t b = \varepsilon \nabla^2 b$

WITH INSULATING: $b(1) = b(-1) = 0$

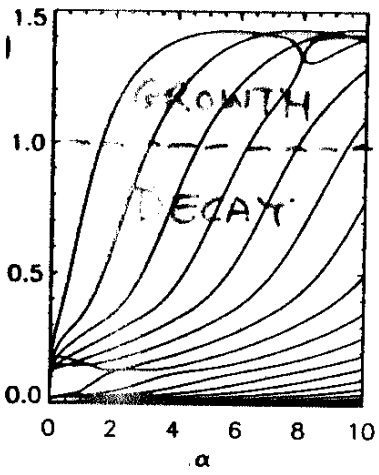
PERFECTLY CONDUCTING: $\partial_x b(1) = \partial_x b(-1) = 0$

OR PERIODIC: $b(x)$ PERIODIC.

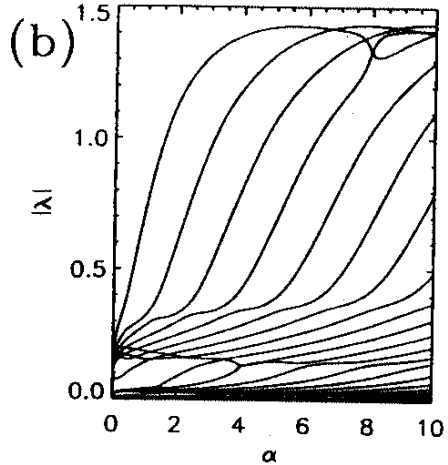
$$H_\varepsilon b(x) = \int_{-1}^1 H_\varepsilon(x,y) b(y) dy$$

EIGENVALUES $|\lambda|$ AGAINST SHEAR PARAMETER α

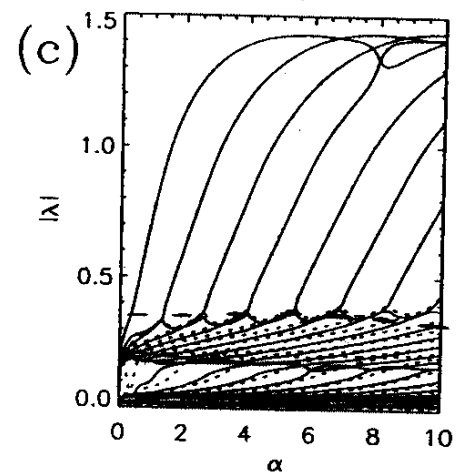
INSULATING



$\epsilon = 10^{-3}$



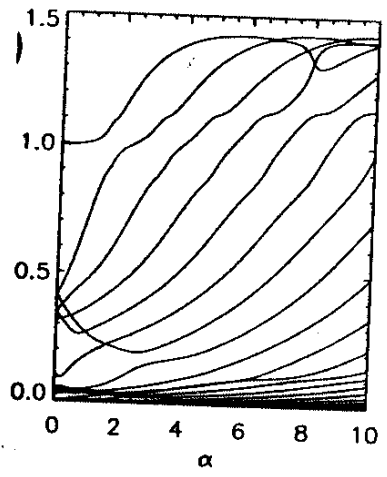
10^{-4}



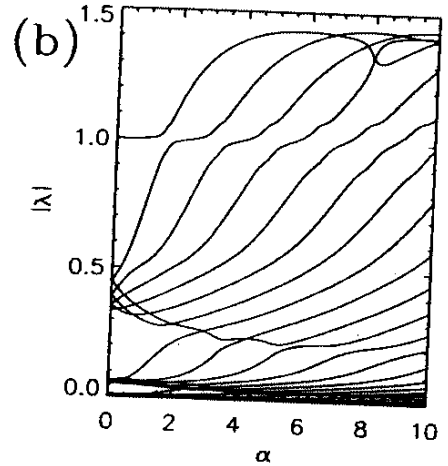
10^{-5}

✓
?

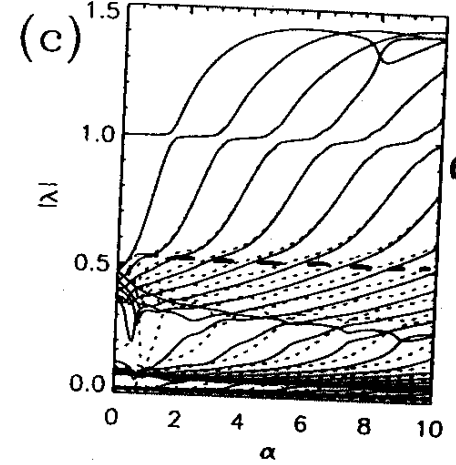
PERFECTLY CONDUCTING



$\epsilon = 10^{-3}$



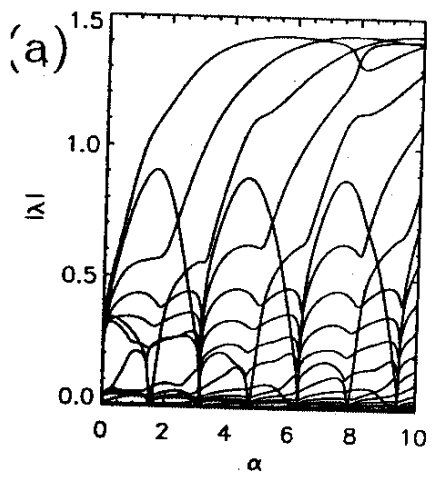
10^{-4}



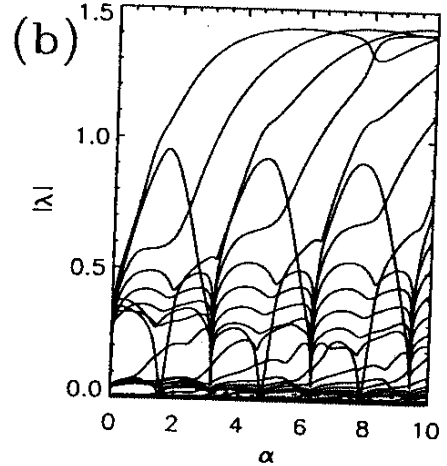
10^{-5}

EXTRA BRANCH
✓
?

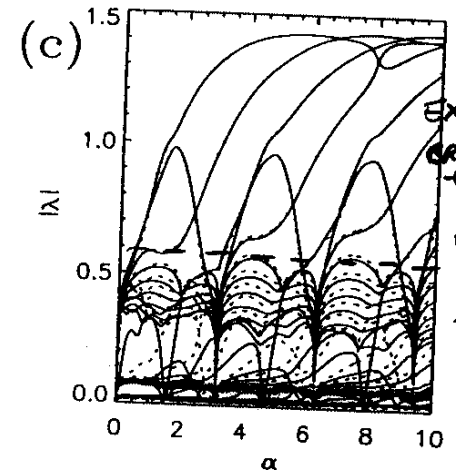
PERIODIC



$\epsilon = 10^{-3}$



10^{-4}



10^{-5}

EXTRA BRANCHES
✓
?

- CLEAR EVIDENCE OF FAST DYNAMO ACTION FOR $\alpha > \pi/2$ WITH ALL BOUNDARY CONDITIONS: PRETTY ROBUST
- SOME DEPENDENCE ON BOUNDARY CONDITIONS: EXTRA BRANCHES, CONVERGENCE, ... ESP. FOR $|\lambda| \leq 1$.

HOW DO WE MAKE SENSE OF THESE RESULTS? AND DEVELOP TOOLS TO ANALYSE THIS & RELATED (E.G. PASSIVE SCALAR) PROBLEMS?

STUDY ADJOINT OPERATOR (IN L^2):

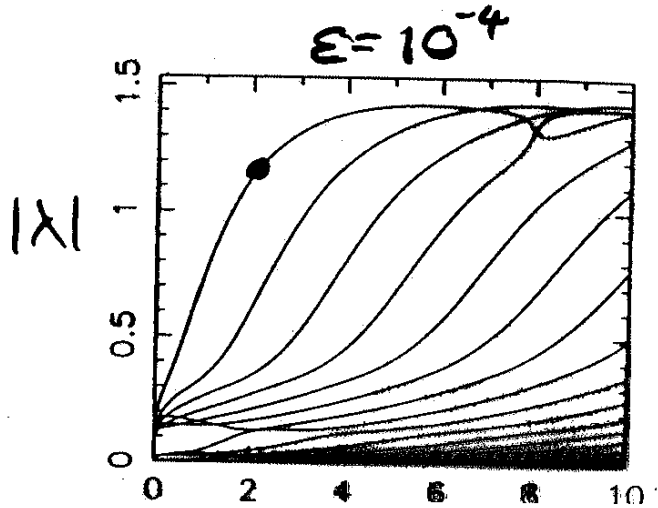
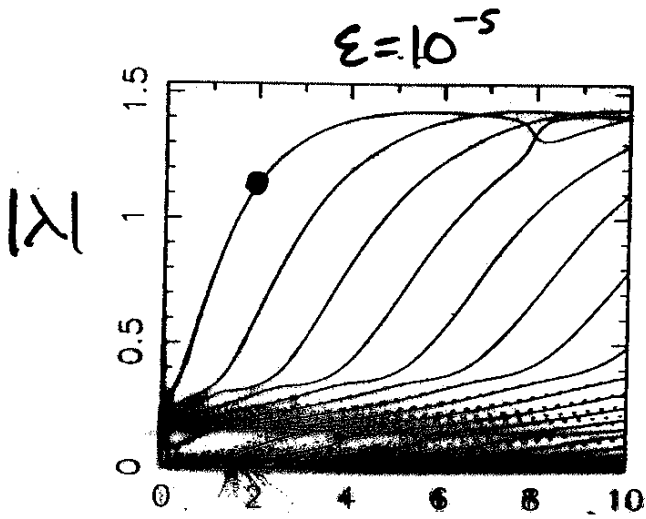
$$T_\varepsilon^* = T^* H_\varepsilon$$

$$T^* c(x) = e^{-i\alpha \frac{1}{2}(x-1)} c\left(\frac{1}{2}(x-1)\right) - e^{-i\alpha \frac{1}{2}(1-x)} c\left(\frac{1}{2}(1-x)\right)$$

Bayly, Childress

THIS STRETCHES OUT STRUCTURE & HAS SMOOTH EIGENFUNCTIONS!

INSULATING $\alpha=2$

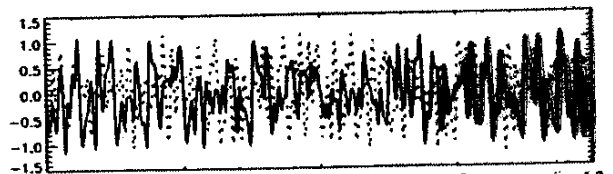
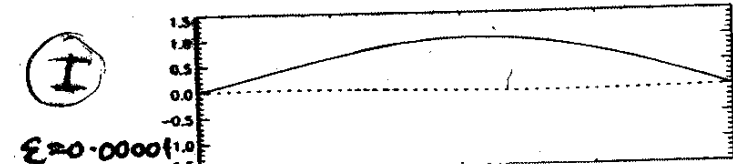
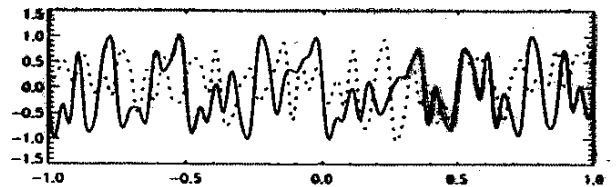
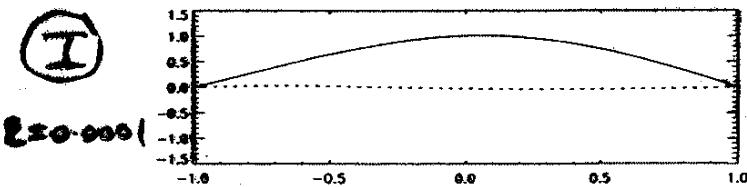
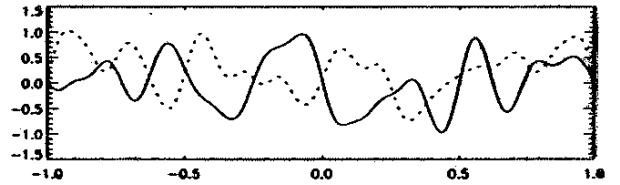
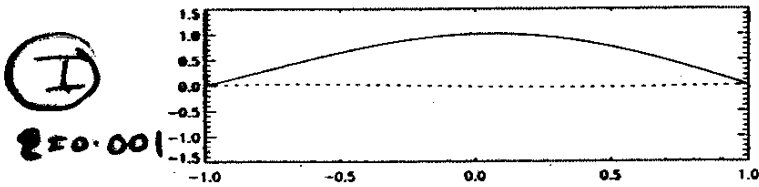
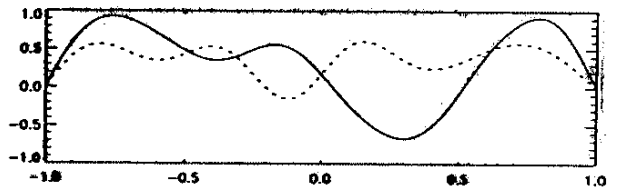
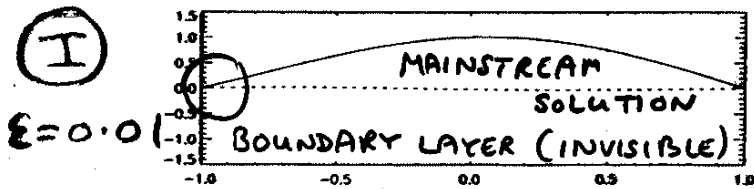


EIGENFUNCTIONS OF

— $Re b$
 $Im b$
 T_E^*

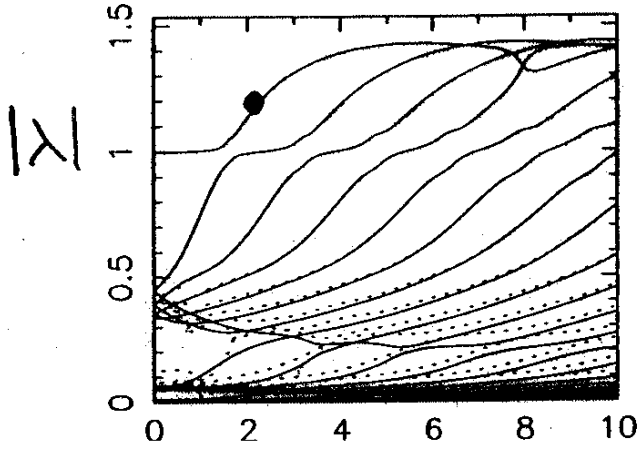
EIGENFUNCTIONS OF

— $Re b$
 $Im b$
 T_E

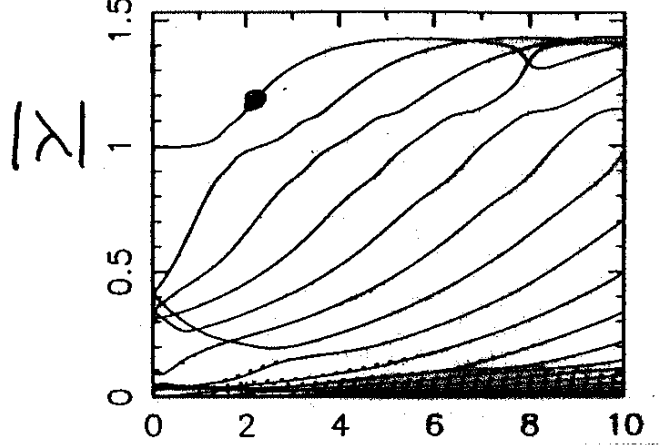


PERFECTLY CONDUCTING
 $\alpha=2$

$\epsilon = 10^{-5}$



$\epsilon = 10^{-4}$

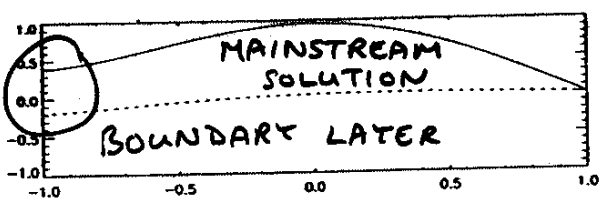


EIGENFUNCTIONS OF

T_ϵ^*

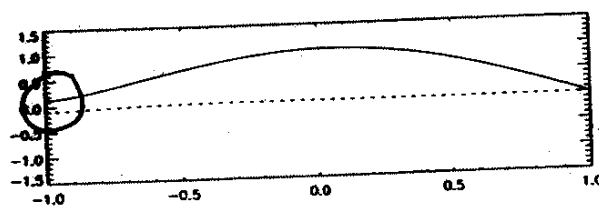
(P)

$\epsilon = 0.01$



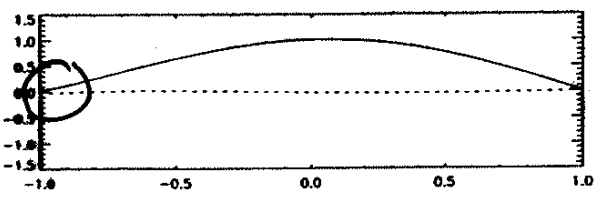
(P)

$\epsilon = 10^{-3}$



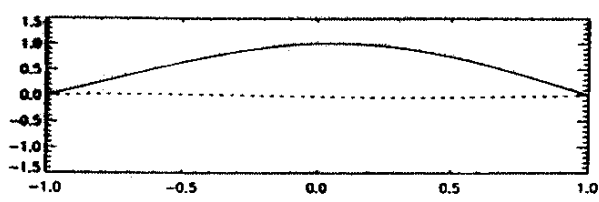
(P)

$\epsilon = 10^{-4}$



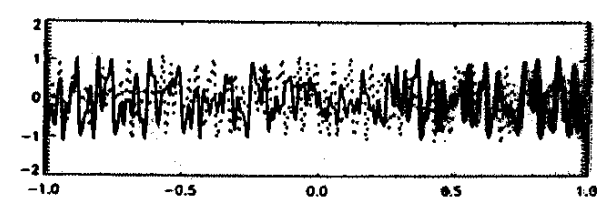
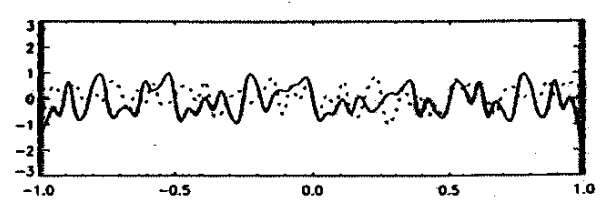
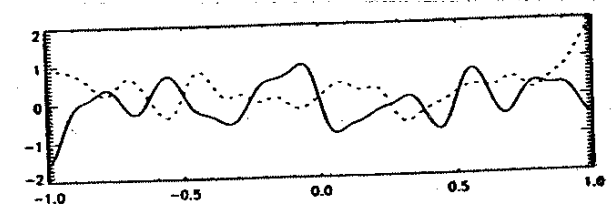
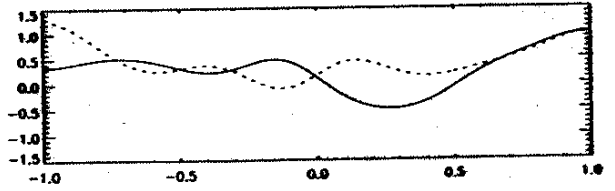
(P)

$\epsilon = 10^{-5}$



EIGENFUNCTIONS OF

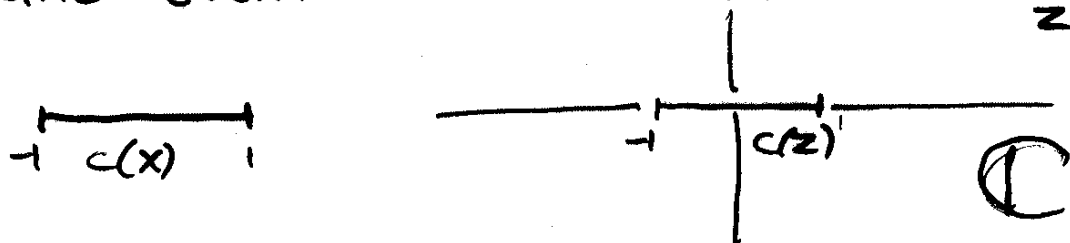
T_ϵ



IN FACT T^* (NO DIFFUSION) HAS

EIGENFUNCTIONS THAT ARE ENTIRE

(ANALYTIC EVERYWHERE IN \mathbb{C})



AND T (NO DIFFUSION) HAS EIGENFUNCTIONS

THAT ARE ANALYTIC EVERYWHERE IN \mathbb{C}

EXCEPT ON $[-1, 1]$.

"HYPERFUNCTIONS"
INCLUDE DISTRIBUTIONS.

E.G. HORIZONTAL BRANCH:

$$T^*: c(z) = e^{i\alpha(z-1)} - e^{i\alpha(1-z)}, \quad \lambda = e^{i\alpha}$$

$$T: b(z) = \frac{e^{-i\alpha}}{z-1} - \frac{e^{i\alpha}}{z+1}, \quad \lambda = e^{i\alpha}$$

CAN CALCULATE EIGENVALUES IN THIS

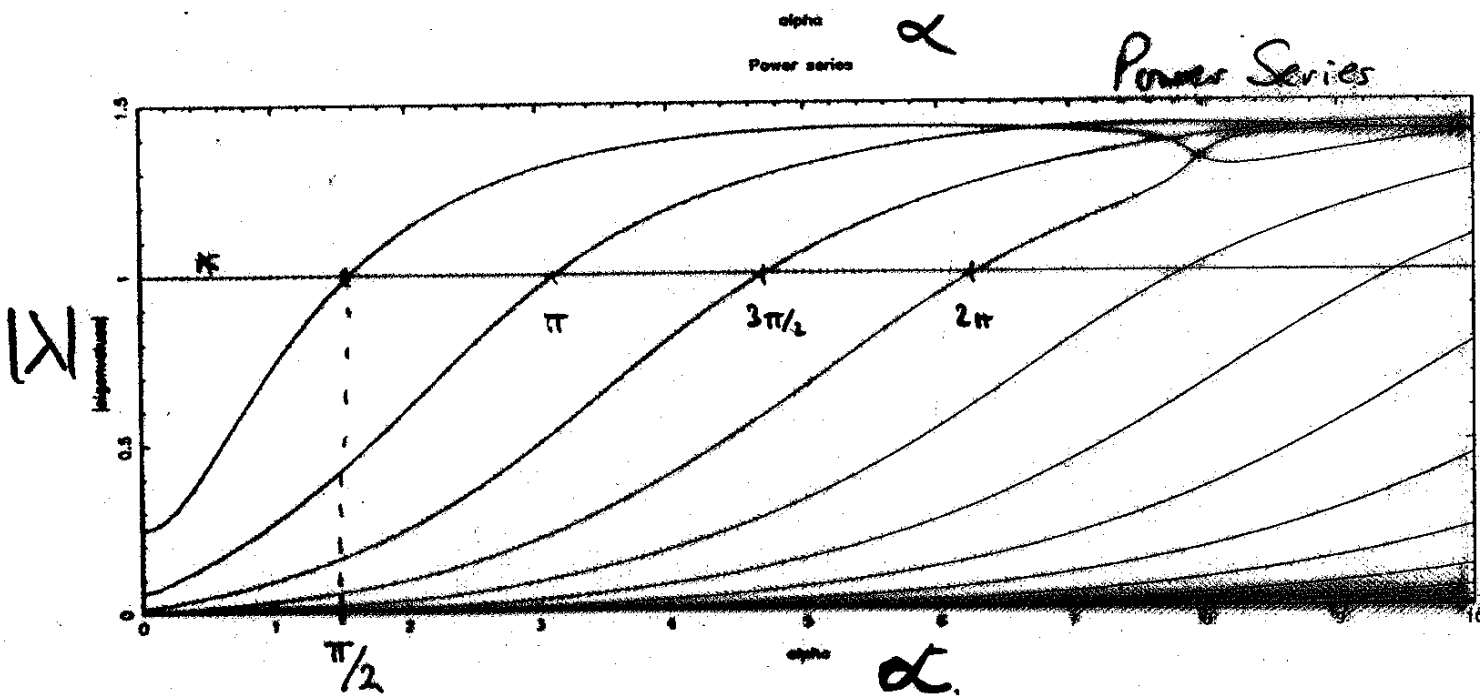
COMPLEX SETTING USING POWER SERIES

EXPANSIONS:

$$c = \sum c_n z^n \quad T^* c = \lambda c$$

$$b = \sum b_n z^{-n-1} \quad T b = \lambda b$$

EIGENVALUES:



ZERO DIFFUSION,

POWER SERIES EXPANSIONS

ENTIRE EIGENFUNCTIONS OF T^*

DISTRIBUTION EIGENFUNCTIONS OF T .

CLOSE RELATION OF EIGENVALUES
 (& EIGENFUNCTIONS OF T^*) WITH ZERO
 DIFFUSION, AND THOSE OF $T_\varepsilon, T_\varepsilon^* \dots$

WRITE:

$$\lambda(\varepsilon) = \lambda(0) + \underbrace{\lambda^{ms} + \lambda^{bl} + \dots}_{\text{LEADING DIFFUSIVE CORRECTIONS}}$$

↑
 EIGENVALUE OF $T_\varepsilon^*, T_\varepsilon$
 FOR WEAK
 DIFFUSION ε
 WITH BOUNDARY
 CONDITIONS

↑
 EIGENVALUE OF T, T^*
 IN-COMPLEX SETTING
 FOR ZERO DIFFUSION

λ^{ms} : EFFECT OF DIFFUSION ON MAINSTREAM COMPONENT
 OF EIGENFUNCTION:

$$\lambda^{ms} = C_1 \varepsilon$$

C_1 CAN BE OBTAINED NUMERICALLY
 FROM EIGENFUNCTIONS.

λ^{bl} : EFFECT OF DIFFUSION IN BOUNDARY LAYERS:

$$\lambda^{bl} = C_2 \varepsilon^2$$

$$q = \frac{1}{2} (n + \log_2 |\lambda|)$$

$q < 0 \Rightarrow$ MODE EVAPORATES.

↑
 STRENGTH OF
 DISCONTINUITY -
 DEPENDS ON BOUNDARY
 CONDITIONS,
 AND EIGENFUNCTION STRUCTURE

C_2 HARD TO CALCULATE.
 NEED TO SOLVE BOUNDARY
 LAYER PROBLEM.....

OPEN QUESTIONS:

- LAMINAR FLOWS: UNDERSTAND/CLASSIFY SEPARATRIX DYNAMOS.
- FAST DYNAMOS IN MORE GENERAL HYPERBOLIC MAPS. (ALSO PASSIVE SCALARS.)
— KOZLOVSKY
- FAST DYNAMOS IN "REALISTIC" FLOWS:
HARD TO ANALYSE; MUCH NUMERICAL EXPLORATION STILL TO BE DONE, E.G. ABC FLOWS.
- DYNAMO SATURATION:

SUPPRESSION OF CHAOS

LIMITING EFFECT OF HELICITY

Vainshtein, Cattaneo, Hughes,
Brandenburg, ...

} SLOWING
OF
TIMESCALE

BUT:

DEPENDENCE ON Re , Rm ,

MODEL DEPENDENCE (DRIVING).

MOSTLY NUMERICAL/CLOSURE.